

Electron heated hybrid JET discharges: experimental results and modeling

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1. Introduction. Tokamak discharges in the hybrid scenario, with a wide area of low magnetic shear (s) and central safety factor (q) close to one, have shown improved MHD stability at high β [1]. Realization of hybrid scenarios with predominant electron heating is a significant step towards ITER-relevant conditions. Such hybrid regimes, with $T_e > T_i$, were first obtained in 2004 in JET; however, they were poorly diagnosed. Recently, these experiments were repeated and extended, with all diagnostics available. After a review of the experiments, this paper presents modeling results, and a comparison with the standard (ion heated) hybrid scenario.

2. Experimental results. Hybrid regimes with $T_e > T_i$ have been obtained in JET by means of Ion Cyclotron Resonance Heating (ICRH, frequency 51.7 MHz) with low-concentration hydrogen minority in deuterium plasmas. Hybrid current profiles were formed by lower hybrid current drive (LHCD) prelude and by adjusting the timing of main heating during the current ramp.

Figure 1 shows the time evolution of reference pulse 68383. Main heating consisted of 9 MW neutral beam injection (P_{nbi}) plus 8.5 MW ICRH, yielding peak temperatures $T_e = 9\text{-}11$ keV and $T_i = 7\text{-}8$ keV, and central toroidal rotation of 30 krad/s, at line averaged density $2.2 \cdot 10^{19} \text{ m}^{-3}$, corresponding to 0.28 times the Greenwald limit (n_G). The confinement regime was H-mode with H89 up to 2 and $\beta_N \simeq 1.5$.

Starting from this reference pulse, other pulses were made by modifying one parameter: (i) a high density pulse with $n_e = 0.57 \cdot n_G$; (ii) a pulse with LH preheat ramped up to 1 MW; (iii) two pulses with P_{nbi} up to 16 MW.

The reference LHCD power was 0.7 MW; with this, tiny sawteeth typically developed after 6 s of main heating. Sawteeth were completely suppressed by ramping the LHCD power to 1 MW; however, in this case a persistent $n = 1$ mode developed [2], see Fig.2. The ELM behaviour depended on n_e and input power: whereas the reference pulse had very frequent, tiny ELMs, the high n_e pulse had small, regular ELMs ($f \sim 100$ Hz), and the high power pulses developed ELM-free periods followed by big ELMs.

3. Interpretative and predictive modeling CRONOS [3] was used both for interpretative and predictive modeling. For the latter, the empirical Bohm/gyroBohm model [4] was used, without taking $E \times B$ shear into account (due to low rotation, the effect of $E \times B$ shear on turbulence reduction is expected to be small). TRANSP [5] was used for calculation of power deposition of ICRH and NBI. We restrict ourselves to the reference pulse, the high density pulse and the high power pulse.

* See Appendix of M.L. Watkins et al., *Fusion Energy 2006 (Proc. 21th Int. Conf. Chengdu) IAEA, paper OV/1-3*

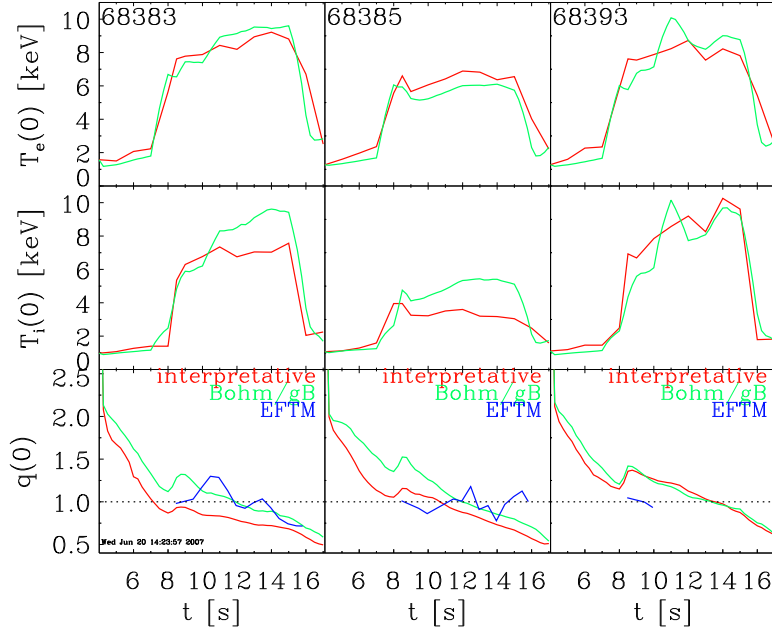


Figure 4: Time traces of central T_e (upper), T_i (middle) and q for the reference discharge (left), high density pulse (middle) and high power pulse (right). Shown are the results of interpretative analysis (red) and predictive modeling with the Bohm/gyroBohm model (green). The lower panels also show the evolution of $q(0)$ from the equilibrium reconstruction using constraints from the MSE measurements (cyan).

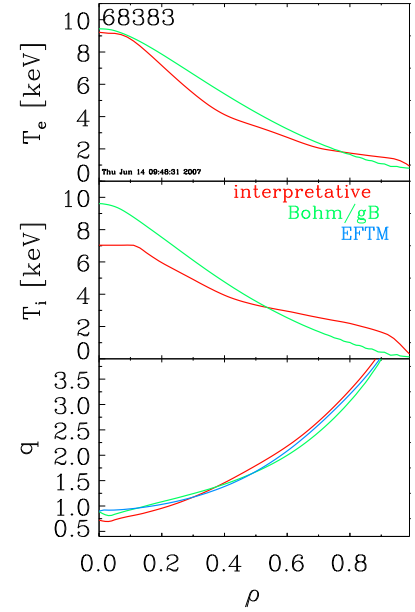


Figure 5: Profiles at $t=14$ s of, from top to bottom, T_e , T_i and q for the same discharges and with the same colour codings as in the previous figure.

4. Comparison with standard hybrid scenario Figures 6, 7 and 8 compare the current reference pulse (68383) with a pulse (62494) in the standard hybrid scenario developed earlier on JET [6], characterized by dominant ion heating (19 MW of NBI and 2 MW of ICRH) and $T_i > T_e$.

Pulse 62494 had much stronger toroidal rotation (fig.6d), hence also a stronger E_r and $\omega_{E \times B}$ shearing rate (fig.7ab; neo-classical v_{pol} was assumed in the calculation). As a very rough measure of the effect of $\omega_{E \times B}$ on turbulence suppression, fig.7c shows $\omega_{E \times B}/\gamma$ where γ is a maximum linear growth rate: around mid radius, it is ~ 0.1 and ~ 0.4 for pulse 68383 and 62494. So the stronger rotation in pulse 62494 is not sufficient to create an ITB; however, the higher value might cause partial turbulence suppression.

Another factor is the ratio T_i/T_e , which is much larger in the standard hybrid pulse (fig.8a). Formulas for the inverse gradient length $R/L_{T_i, \text{crit}}$ for ITG mode turbulence as given by Jenko, Weiland and Romanelli [7,8,9] indicate that $R/L_{T_i, \text{crit}}$ increases with increasing T_i/T_e . In two of the formulas $R/L_{T_i, \text{crit}}$ increases also with s/q ; however, s/q is not much different for the two pulses (fig.8b). Indeed $R/L_{T_i, \text{crit}}$ is clearly higher for pulse 62494 than for 68383 (fig.8c).

In the region roughly between 3.4 and 3.7 m pulse 62494 indeed has much larger $R/L_{T_i, \text{crit}}$, see fig.6c; in this region $R/L_{T_i, \text{crit}}$ for both pulses is in satisfactory agreement with the predictive formulas. The region of high $R/L_{T_i, \text{crit}}$ of pulse 62494 also coincides with the region of enhanced $\omega_{E \times B}$ shearing rate (fig.7c). It is less clear why $R/L_{T_i, \text{crit}}$ for pulse 62494 drops to such a low value inside 3.4 m. One explanation might be the low heat flux in the central area due to the wide power deposition of NBI. The χ_i profiles show a behaviour consistent with $\omega_{E \times B}$ and R/L_{T_i} : outside $\rho \simeq 0.4$ it is clearly lower for pulse 62494 (fig.7d).

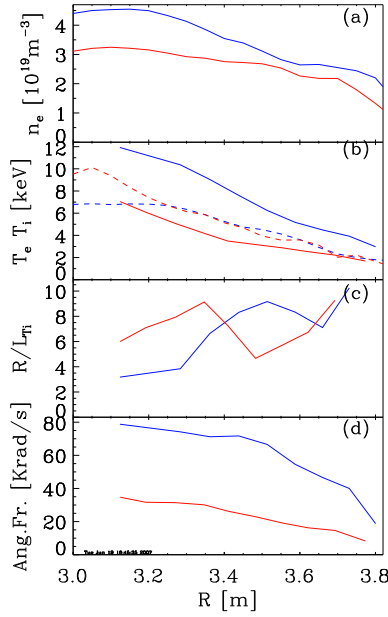


Figure 6: Time averaged profiles during main heating of n_e (a), T_i and T_e (b, full and dashed lines), R/L_{T_i} (c), and toroidal rotation (d), for reference pulse 68383 (red) and the standard ion heated hybrid pulse 62494 (blue).

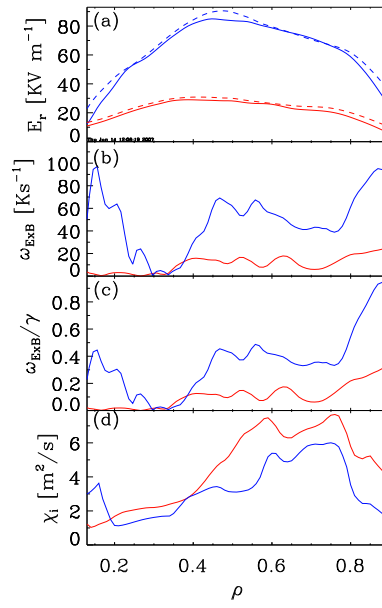


Figure 7: Profiles after 6 s of main heating of E_r (a), $\omega_{E \times B}$ (b), and $\omega_{E \times B}/\gamma$ (c, see text), and χ_i (d), for pulses 68383 (red) and 62494 (blue). In (a) also the contribution to E_r due to the toroidal rotation ($E_{r,\phi}$) is shown (dashed lines).

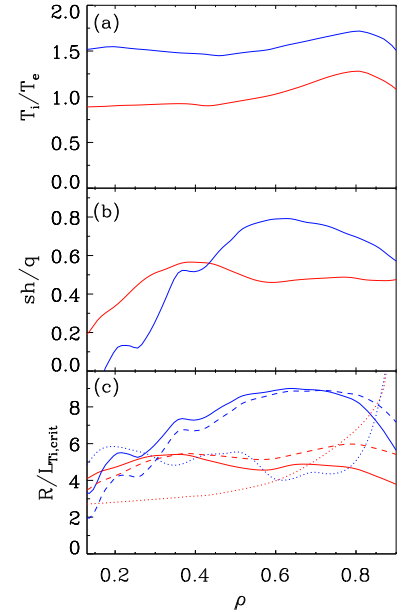


Figure 8: Profiles after 6 s of main heating of T_i/T_e (a), and s/q (b) and $R/L_{T_i, \text{crit}}$ (c) for pulses 68383 (red) and 62494 (blue). Shown in panel (c) are the expressions of Jenko (full), Weiland (dotted) and Romanelli (dashed, see text).

5. Discussion ITER relevant hybrid discharges with $T_e > T_i$ and H-mode factor H89 up to 2 have been obtained in JET. These pulses have been compared with a reference standard (ion heated) hybrid discharge. Both $\omega_{E \times B}$ and T_i/T_e are significantly larger for the latter pulse. In agreement with this, outside 3.4 m the inverse T_i gradient scale length was clearly larger for the ion heated pulse, and χ_i was lower. Apparently, ion thermal confinement is deteriorating when one increases electron heating; this has been seen in other machines as well [10].

The clamping of $q(0)$ to a value $\simeq 1$ is in contrast with simulations which predict $q(0)$ to drop far below 1. Maybe the effect of NTMs on the q profile, as observed in DIII-D, plays a role in keeping $q(0)$ close to 1 in the experiment [11].

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