

## Turbulence intermittency and burst properties in the boundary of TEXTOR tokamak

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### 1. Introduction

Turbulence intermittency at the plasma boundary has long been observed [1-4]. It has been generally found that the intermittent structures are convective in nature and propagate radially carrying a significant fraction of the particle flux [3, 4]. The phenomena have been involved in (i) enhanced erosion of the wall and reduction of the divertor efficiency [2]; (ii) the departure of transport scaling from gyroBohm to Bohm [5]. In recent years, models based on the coherent vorticity [6] and the interchange instability [7, 8] were proposed to account for some of the properties. The bursty transport may also arise due to avalanches in the SOC (self-organized criticality) dynamics [8]. In experiments, although common features have been observed, discrepancies also exist from machine to machine: e. g. (i) in the ADITYA tokamak, the small scale coherent structures are found to play a crucial role in inducing bursty transport [9], while on some other machines the origin of intermittency appears to be different [3]; (ii) in the PISCES linear device, the intermittent transport is linked to certain instabilities and occurs periodically [3]. In this paper, we present the experimental results on the intermittency behavior in TEXTOR by using various analyses to give a comprehensive picture of the intermittent transport.

### 2. Experimental setup

The experiments were performed in the TEXTOR tokamak under the following ohmic discharge conditions:  $R=175\text{cm}$ ,  $a\approx 48\text{cm}$ ,  $B_T=2.33\text{T}$ ,  $I_p=200\text{kA}$ ,  $V_l=1\text{V}$  and  $\langle n_e \rangle = 1.0 \times 10^{19} \text{m}^{-3}$ . The measurements were mainly made in the plasma edge and the scrape-off layer (SOL) by a Langmuir probe array ( $7.5 \times 7.5\text{mm}$ ) consisting of 4 carbon tips with each having 4mm long and 3.5mm in diameter. Two of them were biased as a double probe to record the ion saturation current,  $I_s$ . We simply assume that the fluctuations on  $I_s$  are primarily caused by the density fluctuations. The other two pins were poloidally spaced to measure the poloidal electric field  $E_\theta = (V_{f1} - V_{f2})/d$ , where  $V_{f1}$  and  $V_{f2}$  are the floating potentials detected by the two pins and  $d$  is the distance between them. The turbulence-driven particle flux,  $\Gamma$ , is thus estimated by  $\Gamma = \langle \tilde{n} \tilde{v}_r \rangle = \langle \tilde{n} \tilde{E}_\theta \rangle / B$ . The fluctuation data were digitized at a rate of 500 kHz.

### 3. Results and discussion

The turbulence intermittency can be identified in the raw signals of  $I_s$  by a number of bursts in the “sea of turbulence”. The amplitudes of bursts are much higher than the standard deviation of fluctuations ( $\sigma$ ). For qualifying intermittency, one can calculate the PDF of the fluctuation data ( $x$ ) and compare it to a Gaussian by computing its skewness ( $S = \langle \tilde{x}^3 \rangle / \langle \tilde{x}^2 \rangle^{3/2}$ ) and kurtosis ( $K = \langle \tilde{x}^4 \rangle / \langle \tilde{x}^2 \rangle^2 - 3$ ). For a Gaussian signal,  $S=K=0$ ,

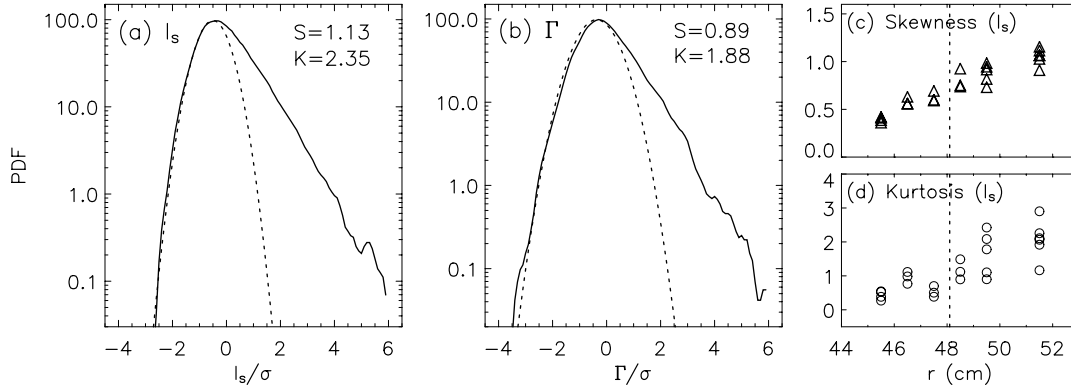
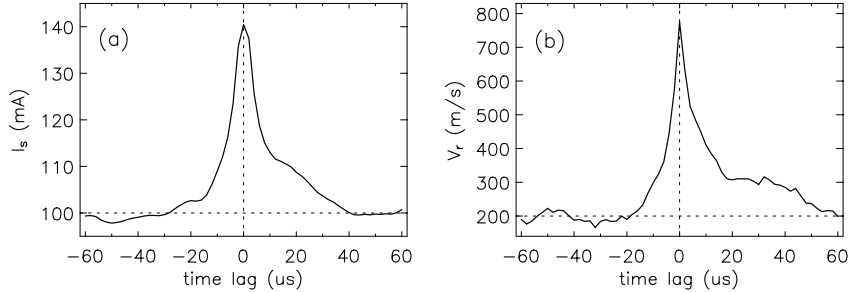


Fig. 1. Semi-log plot of PDFs for (a) the normalized density fluctuations ( $I_s/\sigma$ ) and (b) the turbulent particle flux ( $\Gamma/\sigma$ ) measured in the SOL ( $r=51.5$  cm). Shown in (c) and (d) are the radial profiles of skewness and kurtosis of  $I_s$ , respectively. Vertical dashed lines in (c) and (d) indicate limiter position.

whereas for others the deviation from 0 indicates a higher degree of non-Gaussianity. Plotted in Fig. 1(a) is the PDF of the normalized density fluctuations ( $I_s/\sigma$ ) measured in the SOL. The dotted line shows the best fit by a Gaussian. The PDF skews positively, while the negative fluctuations are close to a Gaussian distribution. The corresponding PDF of particle flux  $\Gamma$  is drawn in Fig. 1(b), which shows similar features to the  $I_s$ . These non-Gaussian PDFs imply a departure of transport properties from pure diffusion to convection. In Figs. 1(c) and (d), the radial dependences of  $S$  and  $K$  of  $I_s$  are plotted. Both of them are enhanced with increasing  $r$ , suggesting increased intermittency from the edge to the SOL. In order to investigate the intermittency structures independently of the background turbulence, the conditional averaging tool [1, 3, 4] is used. In our analysis, the intermittent bursts are selected by their amplitudes by setting a threshold of several times the standard deviation of  $I_s$ . Conditional averaging is achieved by recording 60 data points=120 $\mu$ s around the maximum of each burst, then accumulating and averaging. For auto-conditional average (ACV), the maxima selection and averaging are performed on the same signal, while for cross-conditional average (CCV), the selection is made on one signal and the averaging is done on another. The ACV of density burst detected in the SOL for  $I_{th}/\sigma=3.0$  is plotted in Fig. 2(a). It shows two features: (i) the density bursts consist of merely positive ones, consistent with the PDFs and skewness shown in Fig. 1; (ii) the decay time of the burst is much longer than the rising time. This time-asymmetry points to

relaxation phenomena. These properties are different from those of coherent eddies [6], which usually have a symmetric shape and both positive/negative fluctuations. Using CCV, the intermittent radial velocity  $V_r$  for different levels of  $I_s$  can be determined. The conditional average of  $V_r$  in the SOL for events  $I_{th}/\sigma \geq 2.5$  is drawn in Fig. 2(b). The bursts show a radially outward motion with an average speed  $\approx 450$  m/s. The  $V_r$  shape is also asymmetric,



consistent with the density intermittency.

The contribution from

the intermittent bursts to the total particle flux can be evaluated by calculating their ratio,  $\langle \Gamma_{int} \rangle / \langle \Gamma_{tot} \rangle = \langle I_s E_\theta \rangle_{intermittent} / \langle I_s E_\theta \rangle_{total}$ , where  $\langle \rangle$  denotes an ensemble average. It is found that in the SOL, the bursts with  $I_{th}/\sigma \geq 2.5$  carry a fraction of  $\sim 40\%$  of the total flux. After the bursts selected, the statistics of the waiting time ( $\Delta T$ ) between successive bursts can be obtained. Shown in Fig. 3 are the PDFs( $\Delta T$ ) of  $I_s$  measured at  $r=48.5$ cm for two thresholds ( $I_s/\sigma=2.5$  and 3). It looks that the PDFs have a wide distribution with a maximal  $\Delta T \approx 150 \mu s$ . For large  $\Delta T$  events, the PDF fits well an exponential law for  $I_s/\sigma=2.5$ , whereas for  $I_s/\sigma=3.0$  an extended power-law appears in the tail of PDF. Such a deformation of PDF from small to large burst events seems to be consistent with the SOC hypotheses [8, 10]. To investigate the long-range correlation and self-similar characters of intermittent fluctuations, the auto-power spectrum  $S(f)$  and the Hurst exponent via R/S analysis[11] are calculated for the density fluctuation data. The results for  $I_s$  measured in the SOL are plotted in Fig. 4. The  $S(f)$  displays three distinct decay regions with scaling indices of 0,  $-1$  and  $-1.8$ , respectively, consistent with the sandpile modeling[12]. The  $f^{-1}$  scaling indicates avalanche interactions. The power-law dependence of the spectrum reflects the self-similarity. In the R/S plot, Hurst parameter is fit as  $H=0.71(>0.5)$ , indicating a long-correlation (or self-similar) process. The results are similar to those of potential fluctuations [13], supporting the idea that the SOC dynamics may play a crucial role in the intermittent transport [8].

#### 4. Conclusion

Turbulence intermittency and burst properties at the boundary of TEXTOR have been investigated. The PDFs of density and particle flux fluctuations are both skewed positively with a Gaussian distribution for the negative fluctuations. The intermittency

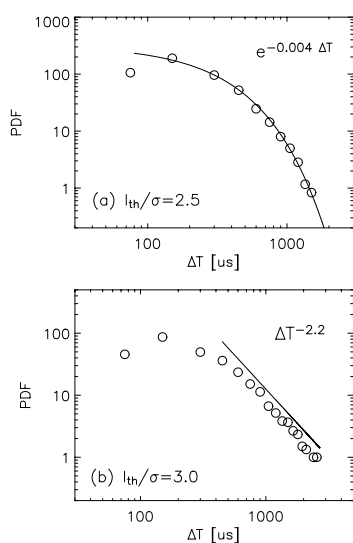


Fig. 3. The PDFs of waiting-time between successive bursts in  $I_s$ , detected at  $r=48.5$  cm for (a)  $I_{th}/\sigma=2.5$  and (b)  $I_{th}/\sigma=3.0$ .

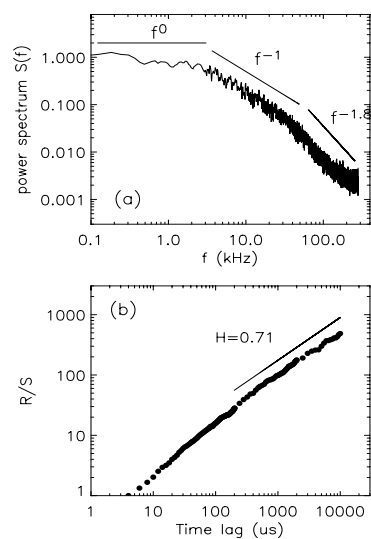


Fig. 4. (a) Power spectrum  $S(f)$  and (b) R/S analysis of density fluctuations measured at  $r=51.5$  cm.

becomes stronger from the plasma edge to the SOL. The averaged shape of bursts is time-asymmetric and composed of only positive ones. In the SOL, the bursts travel radially outwards at a speed of  $\sim 450$  m/s, carrying  $\sim 40\%$  of the total particle flux. The PDF of the waiting-time exhibits a Poisson-distribution for small-size events and a power-law for larger-size ones. The  $S(f)$  and Hurst parameter reveal self-similarity of the intermittent fluctuations. The results indicate that the intermittency in TEXTOR is not dominated by coherent vortices and probably related to avalanche-like transport.

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