

## Study of Resonance Helical Field on the $Z_{\text{eff}}$ and Impurity Radiation in IR-T1 Tokamak

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The RHF in tokamak is an external magnetic field which can improve the plasma confinement. This field is produced by conductors wound externally around the tokamak torus with a given helicity. In IR-T1, RHF is generated by two of helical coils installed outside the vacuum vessel. The pulsed dc RHF configuration ( $\ell = 2,3$ ) has the optimal current and variable time. The magnitude of the weak RHF used in our experiments did not exceed 1% of Bp. The RHF was not effective in high  $Z_{\text{eff}}$  discharges. The RHF could reduce energy and particle losses. The aim that is followed in this article is to understand the effect of RHF on the  $Z_{\text{eff}}$  in addition to influence on the control of plasma. Probably the RHF produces a fixed helical structure with in the plasma which hinders a rotation of the MHD modes and also there with a convective growth of the perturbation [1]. Since the facts of how RHF can affect on plasma control, is not clear, so it seems to be difficult to understand how RHF can impress  $Z_{\text{eff}}$ . To investigate this effect, in this article we will outline the calculation of  $Z_{\text{eff}}$  value through anomaly factor by RHF and without RHF.

The anomaly factor  $\alpha$  is defined as the ratio of the measured plasma resistivity  $\eta_p$  to the theoretical resistivity  $\eta_{11}$  predicted by spitzer, for pure hydrogen plasma [2]. The spitzer

resistivity is  $\eta_{11} = 5.24 \times 10^{-5} \frac{Z \ln \Lambda}{T_e^{3/2}} (\Omega \cdot m)$ , Where Z is the ion atomic mass,  $\ln \Lambda$  is the

coulomb logarithm and  $T_e$  is the electron temperature in eV. For the estimation of the electron temperature it was necessary to perform experimental measurements of plasma current, loop voltage and electron density and to take into account the geometrical parameters of IR-T1

tokamak. The coulomb logarithm is  $\Lambda = 1.5492 \times 10^{13} \frac{T_e^{3/2}}{Z^2 \sqrt{\bar{n}_e}}$ , Where  $\bar{n}_e$  is the average

plasma density in  $m^{-3}$ . The electron temperature is calculated according to the energy balance

equation for a tokamak discharge:  $T_e = \frac{I_p V_\ell T_E}{2k\pi^2 R_r^2 \bar{n}_e}$ , Where  $I_p$  is the plasma current is A,  $V_\ell$

is the loop voltage in V,  $T_E$  is the energy confinement time in s, k is the Boltzman constant, R

and r are the IR-T1 major and minor radii in m. for ohmic heating, plasma which is not too

dense, the energy confinement time is given as  $T_E = 7 \times 10^{-22} \bar{n}_e r R^2 q$ , Where q is the safety

factor which is given as  $q = 5 \times 10^6 \frac{B_T r^2}{I_p R}$ , Where  $B_T$  is the toroidal magnetic field in

tokamak. The values of q,  $T_E$ ,  $\ln \Lambda$ ,  $T_e$ ,  $\eta_{11}$  and  $\eta_p$ , for the IR-T1 tokamak were calculated

from experimental data as bellow:  $I_p = 30kA$ ,  $V_\ell = 1.5V$ ,  $n_e = 1.1 \times 10^{13} cm^{-3}$ ,  $B_T = 0.6 T$

The geometrical parameters of the IR-T1 tokamak are: R= 45 cm, r = 12.5 cm. Thus, the

resulting values are:  $q = 3.47$ ,  $T_E = 0.7 ms$ ,  $T_e = 129 eV$ ,  $\Lambda = 6.9 \times 10^6$ ,  $\ln \Lambda = 15.75$ , For

hydrogen plasma in the IR-T1 tokamak (Z=1) the spitzer resistivity is

$\eta_{11} = 0.56 \times 10^{-6} (\Omega m)$ , For Ohmic input power, the plasma resistivity is given by

$\eta_p = \frac{r^2 V_\ell}{2 R I_p}$ , value of  $\eta_p$  for IR-T1 tokamak is equal with  $0.86 \times 10^{-6}$  and because we have

defined  $Z_{eff}$  as the ratio of the measured plasma resistivity  $\eta_p$  to the theoretical resistivity  $\eta_{11}$ ,

the value of  $Z_{eff}$  for the IR-T1 tokamak is 1.5 . When RHF is on, the experimental and

theoretical results [3] show no Change in the main field components, toroidal. Since the

toroidal component mains constant, we expect that the  $Z_{eff}$  value will be unchanged through

anomaly factor but experimental results show, although  $I_p$  and  $B_T$  value is unchanged but

remarkable change is seen on the value of  $n_e$  and  $T_e$ . To provide these changes in anomaly

factor, there is noticeable decrease at  $Z_{\text{eff}}$ . Keeping previous experimental conditions, following results are given:

$$n_e = 1.6 \times 10^{13} \text{ cm}^{-3}, q = 3.47, T_E = 0.98 \text{ ms}, T_e = 99.5 \text{ eV}$$

$$\Lambda = 3.8 \times 10^6, \ln \Lambda = 15.15, \eta_{11} = 0.86 \times 10^{-6}, \eta_p = 0.86 \times 10^{-6} \quad \text{and} \quad Z_{\text{eff}} = \frac{\eta_p}{\eta_{11}} = 1.075$$

## References

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3. W N-C Sy, J. Phys. A: Math. Gen. 14(1981)2095-2112.