## Study of Runaway Electrons in the Iran Tokamak 1 (IR-T1)

F. Hajakbari<sup>1,2</sup>, A. Hojabri<sup>2</sup> and M. Ghoranneviss<sup>1</sup>

<sup>1</sup>Plasma Physics Research Center, Islamic Azad University - Science and Research Campus,

Tehran, Iran.

<sup>2</sup>Physics group, Islamic Azad University - Karaj Branch, Iran.

In a fully ionized gas electrons with energies higher than certain critical energy are continuously accelerated when an electric field is applied. The resulting high energy electrons are called runaway electrons.<sup>1</sup> In tokamaks runaway electrons are often observed during and after a plasma disruption.<sup>2</sup> This issue is particularly important in the case of disruption generated runaway electrons with high energy, due to the damage first wall materials.<sup>3</sup> For this reason, experimental and theoretical studies are being carried out in order to determine the conditions which can provide a reduction of the effect of runaway electrons during disruptions.<sup>4-6</sup> Runaway electrons are generated when the accelerating electric field acting on the electrons exceeds decelerating effect that such as collisions losses.

The equation describing the motion of a test particle in momentum space were already obtained from the kinetic Fokker-Planck equation for the fast electron distribution. The dynamics of a relativistic electron in a tokamak plasma will be described using the test particle equations:

$$\frac{dP_{II}}{dt} = eE_{II} - \frac{n_e e^4 Ln \Lambda m_e \gamma}{4\pi\varepsilon_o^2} \left( Z_{eff} + 1 + \gamma \right) \frac{P_{II}}{P^3} - F_s \frac{P_{II}}{P}$$
(1)

$$\frac{dP}{dt} = eE_{II} \frac{P_{II}}{P} - \frac{ne^4 Ln\Delta m_e}{4\pi\varepsilon_o^2} \cdot \frac{\gamma^2}{P^2} - F_s \tag{2}$$

Where  $P_{II}$  and P are the electron momentum parallel to the magnetic field and total electron momentum, respectively,  $\gamma$  is the relativistic gamma factor,  $n_e$  is the line-

averaged central electron density, the e is the absolute value of the electron charge,  $m_e$  is the electron mass,  $\varepsilon_o$  is the vacuum permittivity,  $E_{II}$  is the toroidal electric field,  $Z_{eff}$  is the effective ion charge,  $Ln\Lambda$  is the Coulomb logarithm and  $F_s$  is the decelerating force due to the synchrotron radiation losses. Martin-Solis and et al rewrite these equations in normalize form<sup>7</sup>:

$$\frac{dq_{II}}{d\tau} = D - \gamma(\alpha + \gamma) \frac{q_{II}}{q^3} - (F_{gc} + F_{gv} \frac{q_{\perp}^2}{q^4}) \gamma^4 \left(\frac{v}{c}\right)^3 \frac{q_{II}}{q}$$
(3)

$$\frac{dq}{d\tau} = D\frac{q_{II}}{q} - \frac{\gamma^2}{q^2} - \left(F_{gc} + F_{gy}\frac{q_{\perp}^2}{q^4}\right)\gamma^4 \left(\frac{v}{c}\right)^3 \tag{4}$$

Where  $q = P/m_e c$  and  $\tau = v_r t$  with  $v_r = \frac{n_e e^4 L n \Lambda}{4\pi \varepsilon_o^2 m_e^2 c^3}$  and  $D = \frac{E_{II}}{E_R}$  is the normalized

electron field, with  $E_R = \frac{KT_e}{m_e c^2} E_D$ , where  $E_D = \frac{n_e e^3 L n \Lambda}{4\pi \varepsilon_o^2 K T_e}$  is the Dreicer field,

$$\alpha = 1 + Z_{eff}$$
,  $v$  is the electron velocity and  $F_{gc} = F_{gv} \left( \frac{m_e c}{eB_o R} \right)^2$ ,

 $F_{\rm gy} = \frac{2\varepsilon_{\rm o}B_{\rm o}^2Ln\Delta m_e}{3n_e}$  ( R is the plasma major radius and  $B_{\rm o}$  is the toroidal magnetic

filed) are parameters describing the two contributions to the radiation losses coming from the guiding center motion and the electron gyro-motion, respectively. Typical plasma parameters were<sup>8</sup>:

 $n_e \cong 2 \times 10^{13} / cm^3$ ,  $KT_e = 100 eV$ ,  $B_o = 0.7 Tesla$ ,  $Z_{eff} = 2.3$  (therefore  $\alpha = 3.3$ ), R = 45 cm,  $Ln\Lambda = 17$  with this parameters, we calculate  $E_D = \frac{n_e e^3 Ln\Lambda}{4\pi \varepsilon^2 KT} = 88.51 \frac{volt}{m}$ , we use the electric field  $E_{II}$  inferred from the measured

loop voltage ( $E_{II} \cong \frac{V_{loop}}{2\pi R}$ , which, during the flat top, when the current is stationary, can

be assumed to be roughly uniform along the plasma minor radius) therefore

$$E_{II} \cong \frac{V_{loop}}{2\pi R} \cong 0.353 \frac{volt}{m}$$
 and  $E_{R} = \frac{KT_{e}}{m_{e}c^{2}} E_{D} \cong 0.0172$ , thus normalized electric field

is about 
$$D = \frac{E_{II}}{E_R} \cong 20.5$$
, also we calculate  $F_{gc} = F_{gv} \left(\frac{m_e c}{eB_{\circ} R}\right)^2 \cong 2.7 \times 10^{-7}$  and

 $F_{\rm gy} = \frac{2\varepsilon_{\rm o}B_{\rm o}^2Ln\Delta m_e}{3n_e} \cong 6.8\times 10^{-3}$ . The main result of this analysis for normalized

electric field D>1 the critical velocity at which electrons run away is modified by relativistic effects and runaway generation occurs.

## References

<sup>1</sup>H. Dreicer, Phys. Rev. 115, 238 (1959).

<sup>2</sup>R. Yoshino, T. Kondoh, Y. Neyatani, K. Itami, Y. Kawano, and N. Isei, Plasma Phys. Controlled Fusion 39, 313 (1997).

<sup>3</sup>R. Jaspers, N. J. Lopes-Cardozo, and F. C. Schuller, Nucl. Fusion 36, 367 (1996).

<sup>4</sup>J. R. Martin-Solis, B. Esposito, R. Sanchez, and J. D. Alvarez, Phys. Plasma 6, 238 (1999).

<sup>5</sup>D. G. Whyte, T. C. Jernigan, D. A. Humphreys, A. W. Hyatt, C. J. Lasnier, P. B. Parks, T. E. Evans, M. N. Rosenbluth, P. L. Taylor, A. G. Kellman et al., Phys. Rev. Lett. 89, 055001 (2002).

<sup>6</sup>J. R. Martin-Solis, and R. Sanchez, Phys. Plasma 13, 012508 (2006).

<sup>7</sup>J. R. Martin-Solis, J. D. Alvarez, R. Sanchez, and B. Esposito Phys. Plasma 5, 2370 (1998).

<sup>8</sup>M. Ghoranneviss, A. Hojabri, S. Kuhn, Nucl. Fusion, 43, 210 (2003).